# What is boundary operator?

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# Boundary Operators in flat space-time by Extrapolation

One way to address this problem is to start from the Inverse LSZ formula for massless particles

$$G(\lbrace x_i, \delta_i \rbrace) = \prod_i \int d\mu_{p_i} e^{-i\delta_i p_i \cdot x_i} S(\lbrace \delta_i p_i \rbrace), \ \delta_i = \pm 1, \ x_i \to \infty$$
 (1)

where  $d\mu_{p_i}=\frac{d^{D-1}p_i}{(2\pi)^{D-2}2\omega_{p_i}}$  is the Lorentz invariant integration measure on the cone and  $p^0=\omega_p=|\vec{p}|$ . (Minwalla and collaborators, 2311.03443[hep-th])

► *G* is the bulk time-ordered Green's function and *S* is the S-matrix.

### Celestial Amplitudes

Massless Celestial amplitudes are defined as (Pasterski-Strominger-Shao)

$$\mathcal{A}(\{z_i,\bar{z}_i,\Delta_i,\delta_i\}) = \prod_{i=1}^n \int_0^\infty d\omega_i \omega_i^{\Delta_i - 1} S(\{\delta_i \omega_i, z_i, \bar{z}_i\})$$
(2)

where massless momentum is parametrized as

$$p^{\mu} = \omega q^{\mu}(z,\bar{z}) = \omega \left(1 + z\bar{z}, z + \bar{z}, -i(z - \bar{z}), 1 - z\bar{z}\right), \ \omega \ge 0 \tag{3}$$

► This suggests that at least the (naive) Celestial boundary operators are

$$a(z,\bar{z},\Delta,\delta) = \int_0^\infty d\omega \omega^{\Delta-1} a(\delta\omega,z,\bar{z}) \tag{4}$$

where  $a(\delta\omega, z, \bar{z})$  is the (annihilation) creation operator.

## Boundary Operators from inverse LSZ

▶ One can play with the Inverse LSZ and get

$$G(\{x_{i} \to \infty, \delta_{i}\})$$

$$= \prod_{i} \frac{1}{(2\pi)^{2}} \int \frac{d\Delta_{i}}{2\pi} \Gamma(\Delta_{i}) \frac{(-i\delta_{i})^{\Delta_{i}}}{R_{i}^{\Delta_{i}}} \int \frac{d^{2}w_{i}}{|z_{i} - w_{i}|^{2\Delta_{i}}} \mathcal{A}(\{w_{i}, \bar{w}_{i}, 2 - \Delta_{i}, \delta_{i}\})$$
+Contact Terms
$$= \prod_{i} \frac{1}{4\pi} \int \frac{d\Delta_{i}}{2\pi} \Gamma(1 - \Delta_{i}) \frac{(-i\delta_{i})^{\Delta_{i}}}{R_{i}^{\Delta_{i}}} \tilde{\mathcal{A}}(\{z_{i}, \bar{z}_{i}, \Delta_{i}, \delta_{i}\})$$
+Contact Terms,
$$R_{i} \to \infty$$
(Banerjee, 2406.06690 [hep-th], JHEP)

 $ightharpoonup \tilde{\mathcal{A}}(\{z_i, \bar{z}_i, \Delta_i, \delta_i\})$  is the Shadow Celestial Amplitude.

Schematically

$$G(\lbrace x_i \to \infty, \delta_i \rbrace) \sim \sum_{\lbrace \Delta_i \rbrace} \prod_i R_i^{-\Delta_i} \tilde{\mathcal{A}} \left( \lbrace z_i, \bar{z}_i, \Delta_i, \delta_i \rbrace \right) \tag{6}$$

► So the boundary correlation function obtained from Inverse LSZ is

$$\tilde{\mathcal{A}}(\{z_i,\bar{z}_i,\Delta_i,\delta_i\}) = \prod_i \int d^2w_i \frac{1}{|z_i - w_i|^{2\Delta_i}} \mathcal{A}(\{w_i,\bar{w}_i,2-\Delta_i,\delta_i\})$$
 (7)

► So, Boundary operator according to Inverse LSZ is

$$\tilde{a}(z,\bar{z},\Delta,\delta) = \int d^2w \frac{1}{|z-w|^{2\Delta}} a(w,\bar{w},2-\Delta,\delta)$$
 (8)

#### Good News!

► Shadow Celestial amplitudes are non-distributional.

Stress tensor itself is the shadow of the subleading soft graviton.

 $ightharpoonup w_{1+\infty}$  currents are also light transform (half-shadow) of the soft graviton operators.

One is tempted to say that: Shadow Celestial CFT<sub>2</sub> is the true boundary theory.

➤ So the boundary theory is an ordinary CFT<sub>2</sub> with non-distributional correlation functions.

► More evidence from AdS<sub>3</sub>-CFT<sub>2</sub> to follow.